

Chiral extrapolations with small sea quark mass data in two-flavor lattice QCD

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3. Results(Chiral extrapolations)
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1. Introduction

Features of lattice QCD

[Advantage]

- Non-perturbative calculations

[Disadvantage]

- Statistical and systematic errors
 - Continuum extrapolation($a \rightarrow 0$)
→ improvement of actions and operators
 - Renormalization factor
→ non-perturbative renormalization method
 - Finite size effect(studied well)
 - Chiral extrapolation

[Chiral extrapolation]

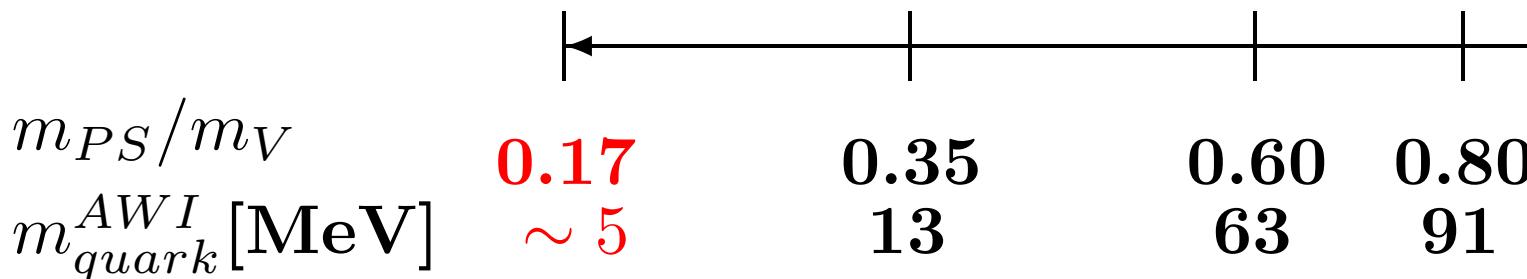
In lattice QCD, sea quark masses

(in two-flavor cases, $m_{ud} \equiv (m_u + m_d)/2$)

are heavier than those in our real world
because of computational costs

→ lattice data must be extrapolated
to the physical point.

lattice $m_{PS}/m_V \rightarrow$ real $m_\pi/m_\rho = 0.17$



- Long extrapolations may involve sizable systematic errors
→ more realistic simulations are desirable

2. Simulations

$N_f = 2$ full QCD simulations

[Action and run parameters]

- Small sea quark masses ($m_{PS}/m_V = 0.60 - 0.35$, $m_{quark}^{AWI} = 63 - 13$ MeV)
- Coarse lattices ($a = 0.2$ [fm]; $L = 2.4, 3.2$ [fm])
- RG gauge + tadpole-improved Clover quark
- HMC with the simple leapfrog scheme
- Measurements : every 5 trajectory (1 conf.)

[Statistics]

Combined with our previous work, CP-PACS,2002

$[12^3 \times 24]$

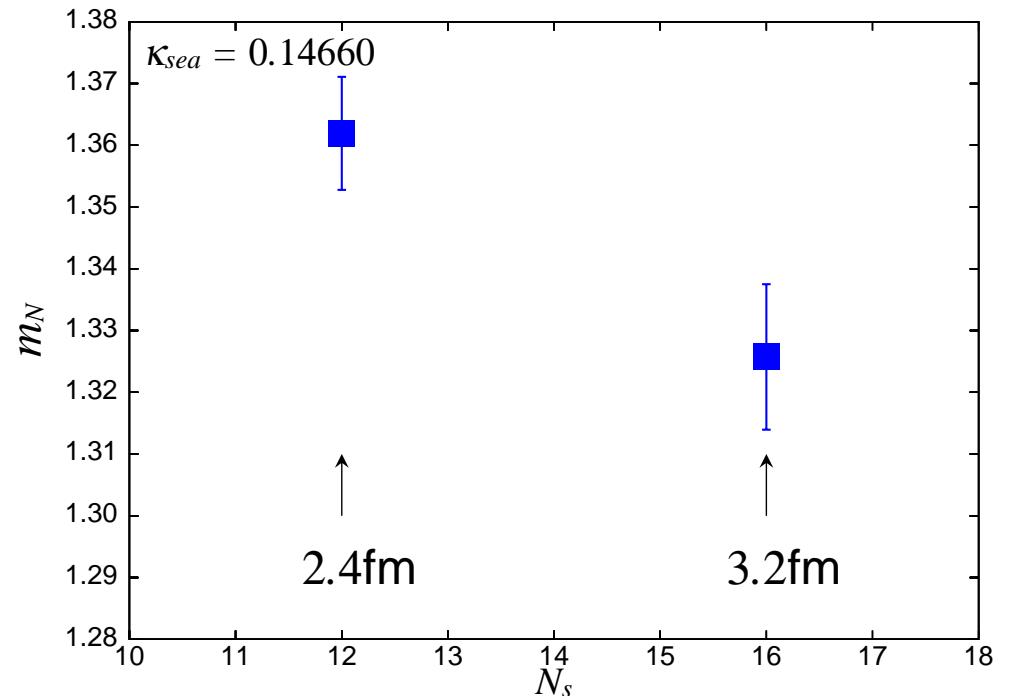
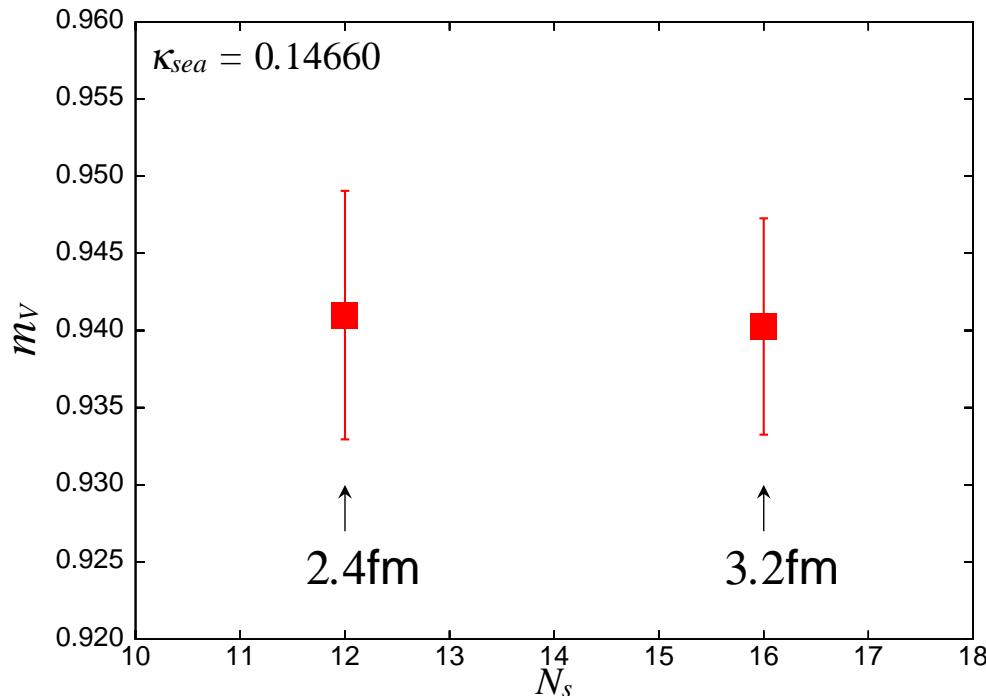
κ_{sea}	m_{PS}/m_V	# traj.
0.14090	0.80	6250
0.14300	0.75	5000
0.14450	0.70	7000
0.14585	0.60	4000
0.14640	0.55	5250
0.14660	0.50	4000
0.14705	0.40	4000
0.14720	0.35	1400

3. Results

[Finite size effect]

$L = 2.4\text{fm}$ vs $L = 3.2\text{fm}$

- No finite size effects in meson quantities ($m_{PS}, m_V, m_{quark}^{AWI}$)
- 1 – 3% ($0.8 – 2.3\sigma$) decreases in baryon quantities (m_N, m_Δ)
→ we concentrate on the meson sector



[Chiral extrapolation]

Parameterize lattice data O with m_{quark}
for chiral extrapolations

$$O = f(m_{quark})$$

Choice of the extrapolation function f

- polynomial (conventional method)
- Chiral perturbation (ChPT) formulae
- Wilson ChPT (WChPT) formulae

Polynomial (conventional) extrapolation

[Vector meson mass]

We parameterize vector meson masses m_V
as a function of pseudoscalar meson masses m_{PS}

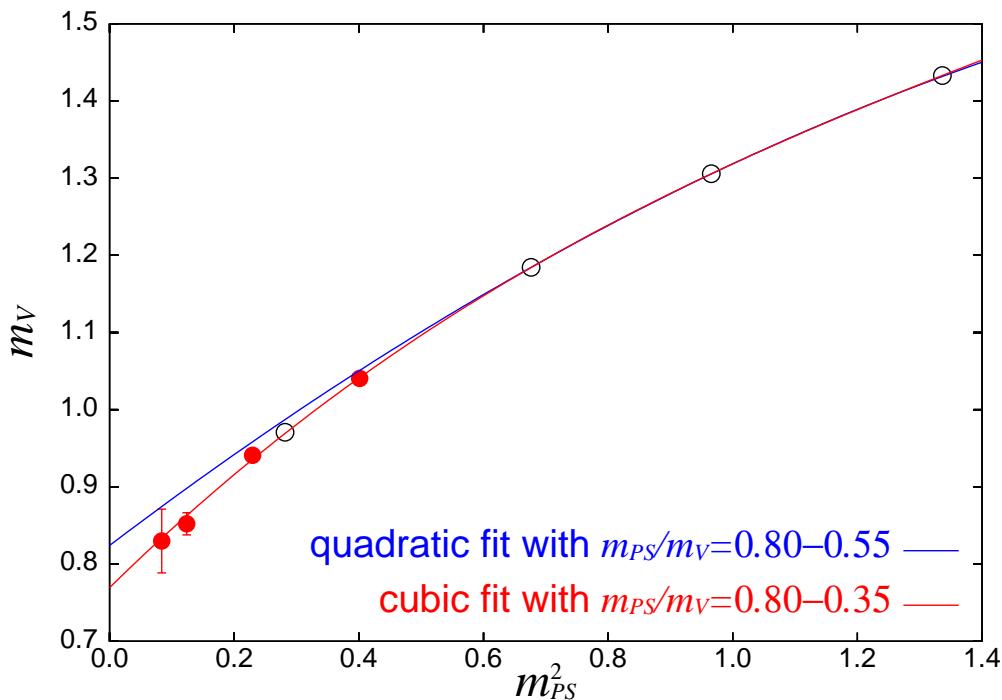
$$m_V = A + Bm_{PS}^2 + Dm_{PS}^4 + Fm_{PS}^6 \text{ for } m_{PS}/m_V = 0.80 - 0.35$$

$$m_V = A + Bm_{PS}^2 + Dm_{PS}^4 \quad \text{for } m_{PS}/m_V = 0.80 - 0.55$$

(previous fit) **CP-PACS,2002**

- Check the validity of the previous quadratic fit from $m_{PS}/m_V = 0.80 - 0.55$ by comparing with new small sea quark mass data $m_{PS}/m_V = 0.50 - 0.35$
- We use m_{PS}^2 instead of m_{quark} for comparison with the previous work (we parameterize m_{PS}^2 as a function of m_{quark} later)

- Systematic deviations from the previous fit are observed in small sea quark mass region
 - Sea quark mass dependence was underestimated by the previous quadratic extrapolation from $m_{PS}/m_V = 0.80 - 0.55$
 - The deviation in the chiral limit is 7%(3.6σ)
 - sizable systematic error!



[Pseudoscalar meson mass]

We parameterize pseudoscalar meson masses m_{PS} as a function of VWI quark masses m_{quark}^{VWI}

$$m_{PS}^2 = B m_{quark}^{VWI} + C(m_{quark}^{VWI})^2 + D(m_{quark}^{VWI})^3 + E(m_{quark}^{VWI})^4$$

for $m_{PS}/m_V = 0.80 - 0.35$

$$m_{PS}^2 = B m_{quark}^{VWI} + C(m_{quark}^{VWI})^2$$

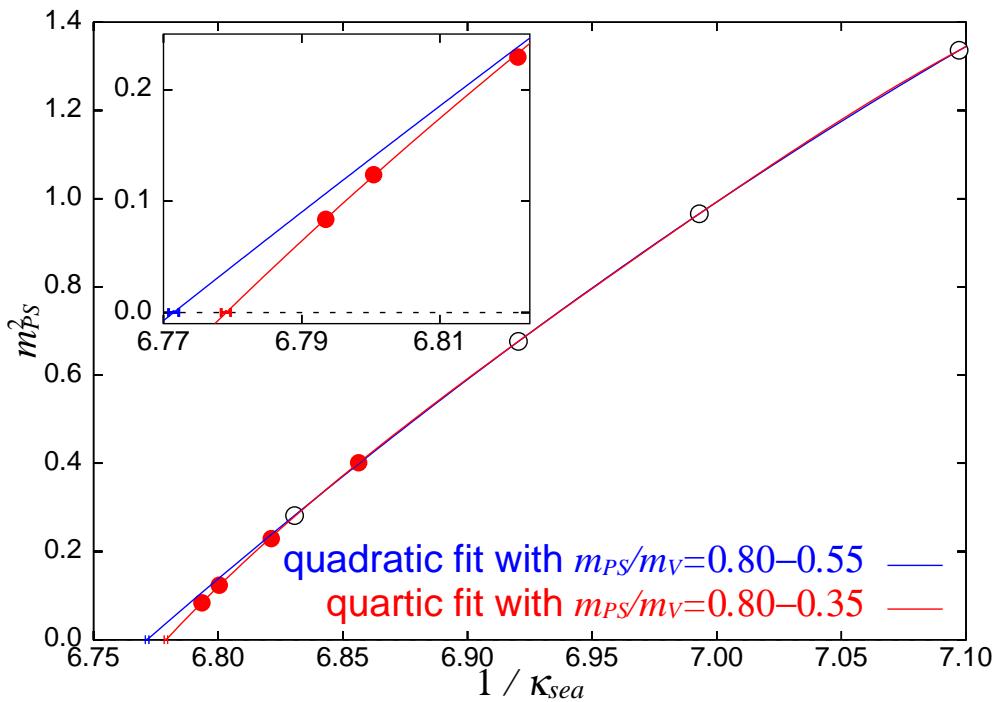
for $m_{PS}/m_V = 0.80 - 0.55$ (previous fit) **CP-PACS,2002**

where

$$m_{quark}^{VWI} = \frac{1}{2} \left(\frac{1}{\kappa} - \frac{1}{\kappa_c} \right)$$

- Up to $(m_{quark}^{VWI})^4$ terms are needed to describe our whole data of $m_{PS}/m_V = 0.80 - 0.35$ with a reasonable $\chi^2/dof = 1.4$

- Large deviations from the previous fit are also observed in small sea quark mass region (10σ difference in κ_c) → sizable systematic error!
- The quadratic fit form can not describe our data of $m_{PS}/m_V = 0.80 - 0.35 (\chi^2/dof = 10)$
cf. $m_{PS}/m_V = 0.80 - 0.55 (\chi^2/dof = 0.1)$

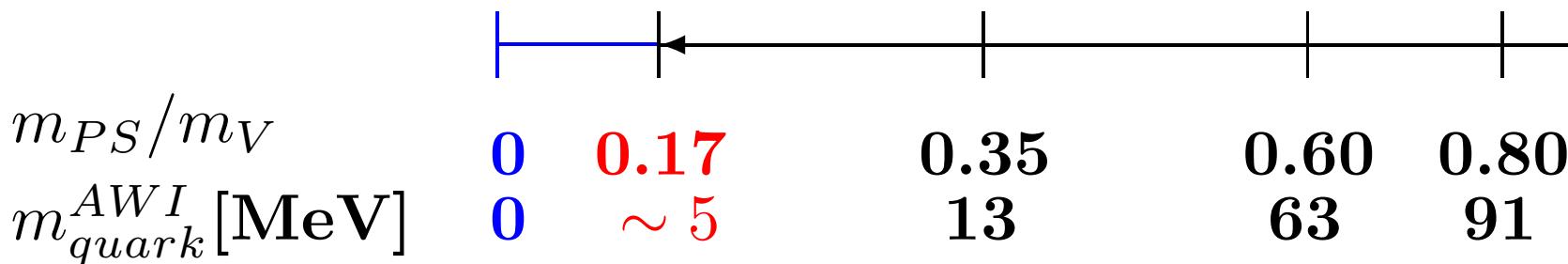


Chiral extrapolations based on ChPT

- In the case of polynomial extrapolations, more higher order terms may be needed if smaller sea quark mass data are available
→ ChPT may give a guide for extrapolations, which is expected to describe the sea quark mass dependence around the chiral limit.

ChPT : a low energy effective theory of QCD

Weinberg,1979;Gasser and Leutwyler,1984,1985

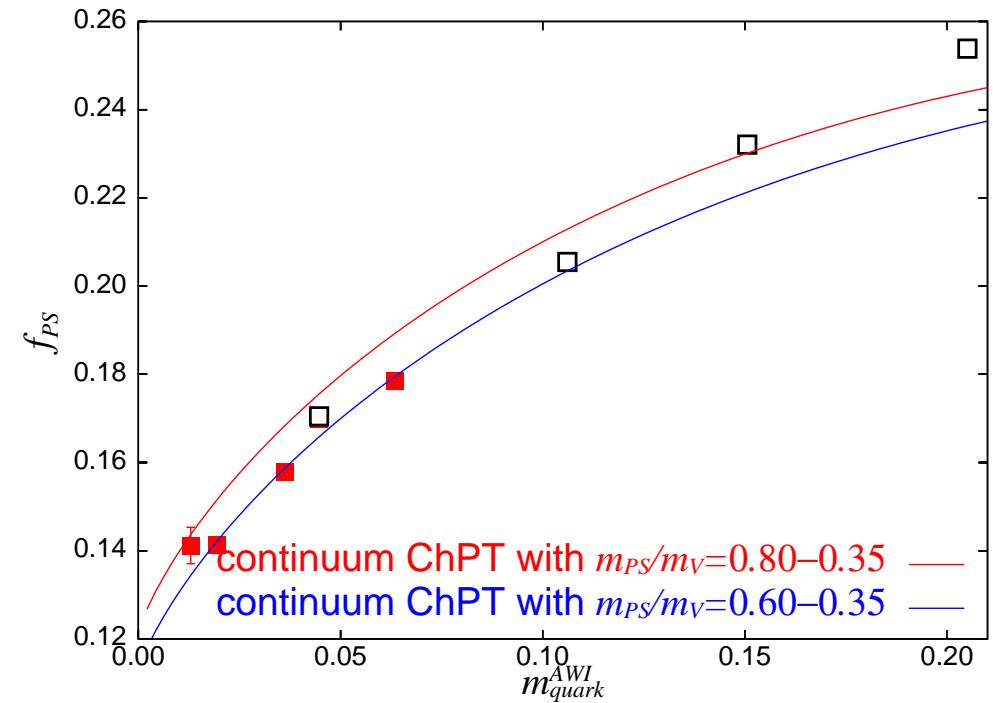
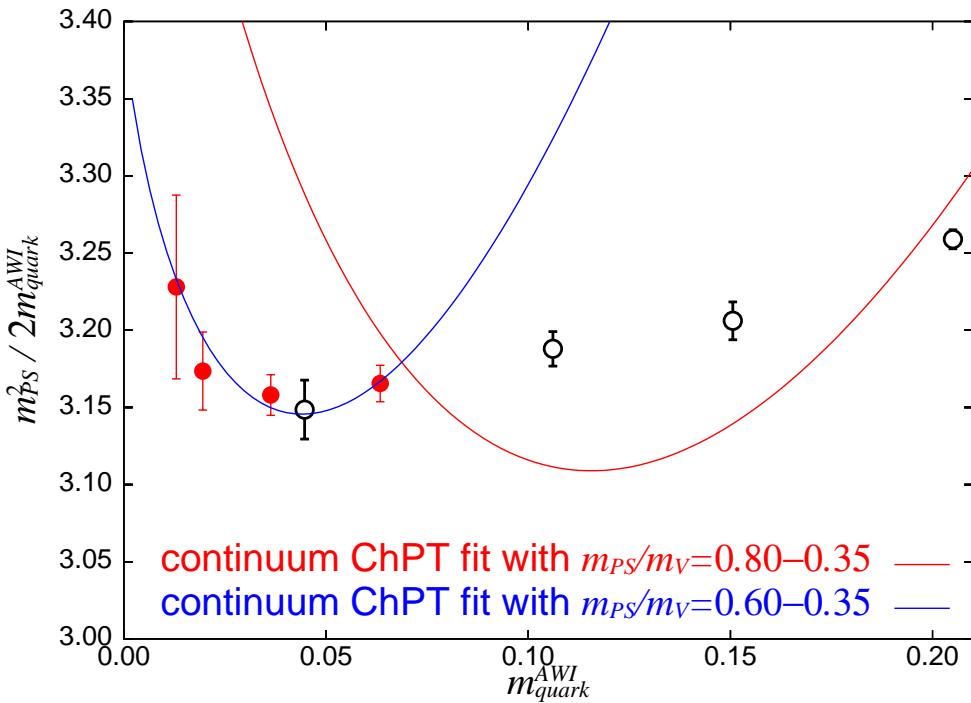


[Simultaneous fit to m_{PS}^2 and f_{PS} with one-loop ChPT]

$$m_{PS}^2 = 2m_{quark}B_0 \left(1 + \frac{m_{quark}B_0}{(4\pi f)^2} \log \frac{2m_{quark}B_0}{\Lambda_3^2} \right)$$

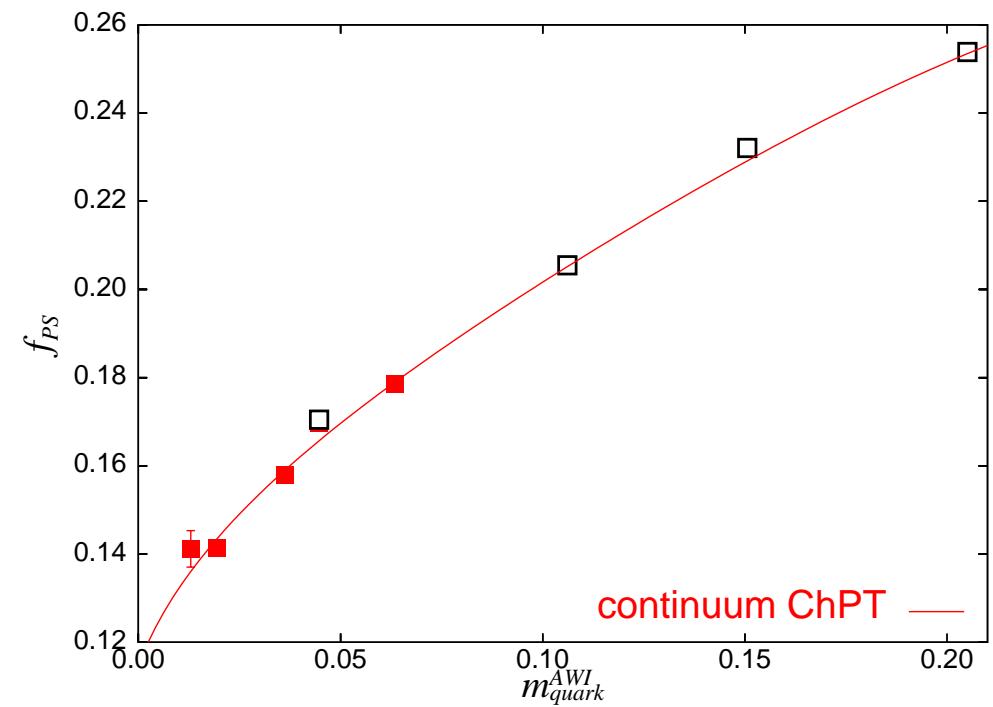
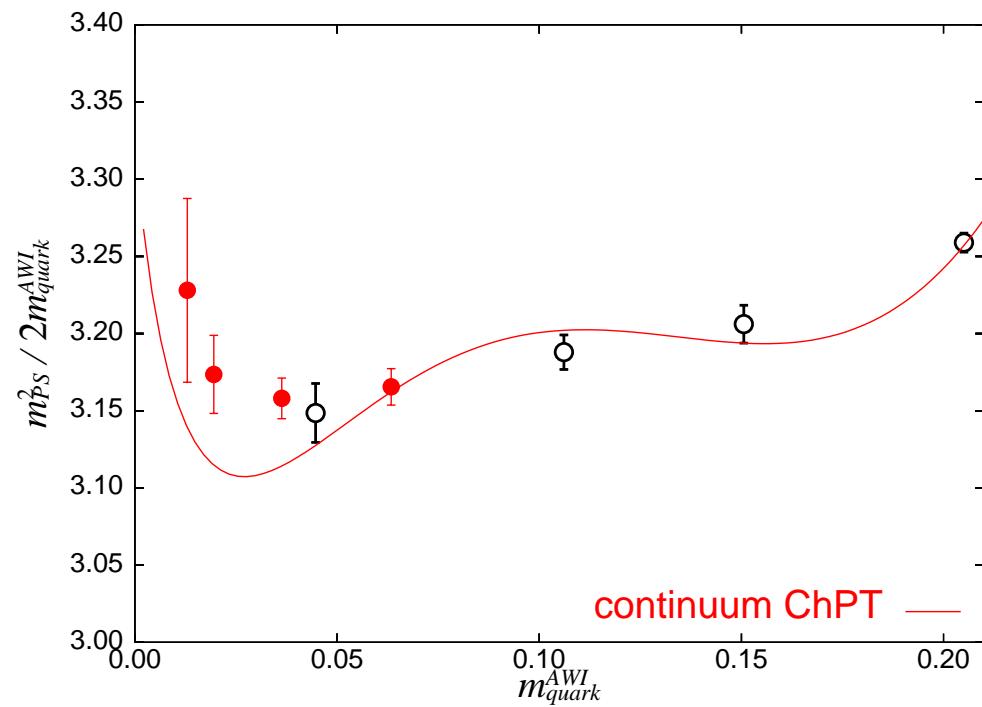
$$f_{PS} = f \left(1 - \frac{2m_{quark}B_0}{(4\pi f)^2} \log \frac{2m_{quark}B_0}{\Lambda_4^2} \right)$$

- ChPT fails to describe our data of $m_{PS}/m_V = 0.80 - 0.35$
- To get a reasonable $\chi^2/dof \sim 4$, we must drop the heavy data
→ ChPT log may appear only below $m_{PS}/m_V \sim 0.4$



[Simultaneous two-loop ChPT fit to m_{PS}^2 and f_{PS}]

- ChPT fails to describe our data of $m_{PS}/m_V = 0.80 - 0.35$ (f_{PS} is OK but mass is not)
- Dropping the heavy data gives a similar wavy fit curve



[Reasons of failure of the continuum ChPT formulae]

- Sea quark masses are still too large for the one-loop and two-loop formulae
→ ChPT log may appear only below $m_{PS}/m_V \sim 0.4$
- Chiral symmetry is explicitly broken for the Wilson-type fermions
→ Modify ChPT for Wilson-type fermions (WChPT)
Sharpe and Singleton, 1998; Rupak and Shores, 2002;
O.Baer et.al., 2003; S.Aoki, 2003

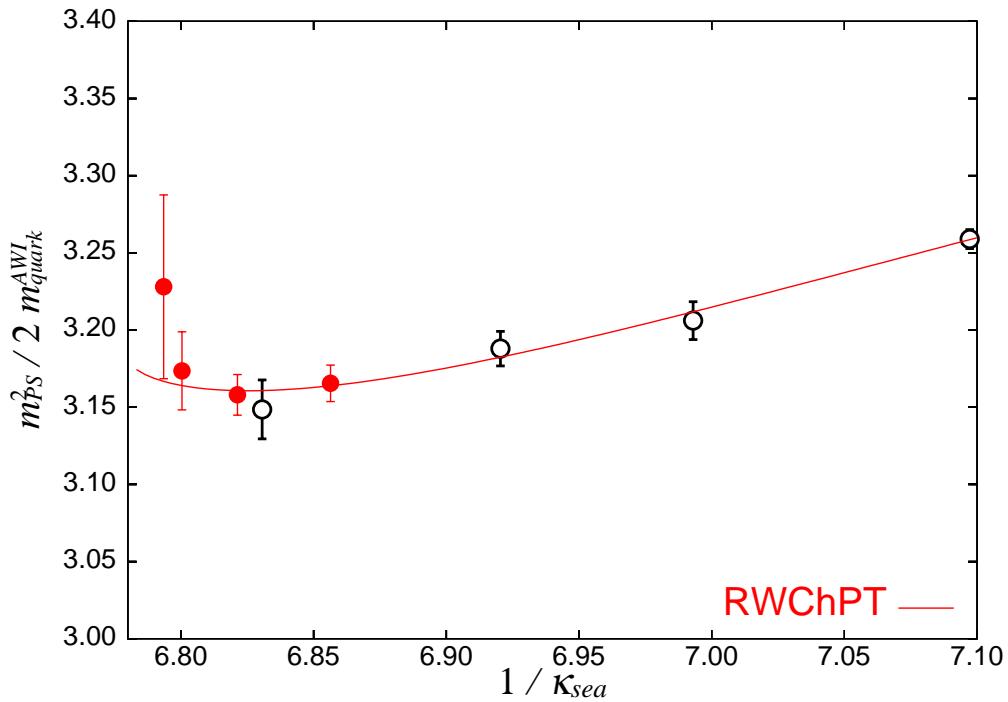
$$\begin{aligned} \mathcal{L}_{WChPT} &= \frac{f^2}{4} \left(1 + \textcolor{blue}{c}_0 (S^0 - 1) \right) \text{Tr} (\partial_\mu U \partial^\mu U) + \textcolor{blue}{c}_1 S^0 + \textcolor{blue}{c}_2 (S^0)^2 \\ S_0 &\equiv \text{Tr}(U + U^\dagger), \quad U \equiv \exp \left(i \sum_{a=1}^3 \frac{\pi_a \sigma^a}{f} \right) \in SU(N_f = 2) \\ c_0 &= W_0 \textcolor{blue}{a}, \\ c_1 &= W_1 \textcolor{blue}{a} + B_1 \textcolor{red}{m}_{quark}, \\ c_2 &= W_2 \textcolor{blue}{a}^2 + V_2 (\textcolor{red}{m}_{quark} \textcolor{blue}{a}) + O(\textcolor{red}{m}_{quark}^2). \end{aligned}$$

where $W = V = 0$ recovers the continuum ChPT Lagrangian

[Resummed WChPT formulae at one-loop] S.Aoki,2003

$$m_{PS}^2 = Am_{quark}^{VWI} \left(-\log \left(\frac{m_{quark}^{VWI}}{\Lambda_0} \right) \right)^{\omega_0} \left(1 + \omega_1^{PS} m_{quark}^{VWI} \log \left(\frac{Am_{quark}^{VWI}}{\Lambda_3^2} \right) \right)$$

$$m_{quark}^{AWI} = m_{quark}^{VWI} \left(-\log \left(\frac{m_{quark}^{VWI}}{\Lambda_0} \right) \right)^{\omega_0} \left(1 + \omega_1^{AWI} m_{quark}^{VWI} \log \left(\frac{Am_{quark}^{VWI}}{\Lambda_{3,AWI}^2} \right) \right)$$



[WChPT fit with our all data of $m_{PS}/m_V = 0.80 - 0.35$]

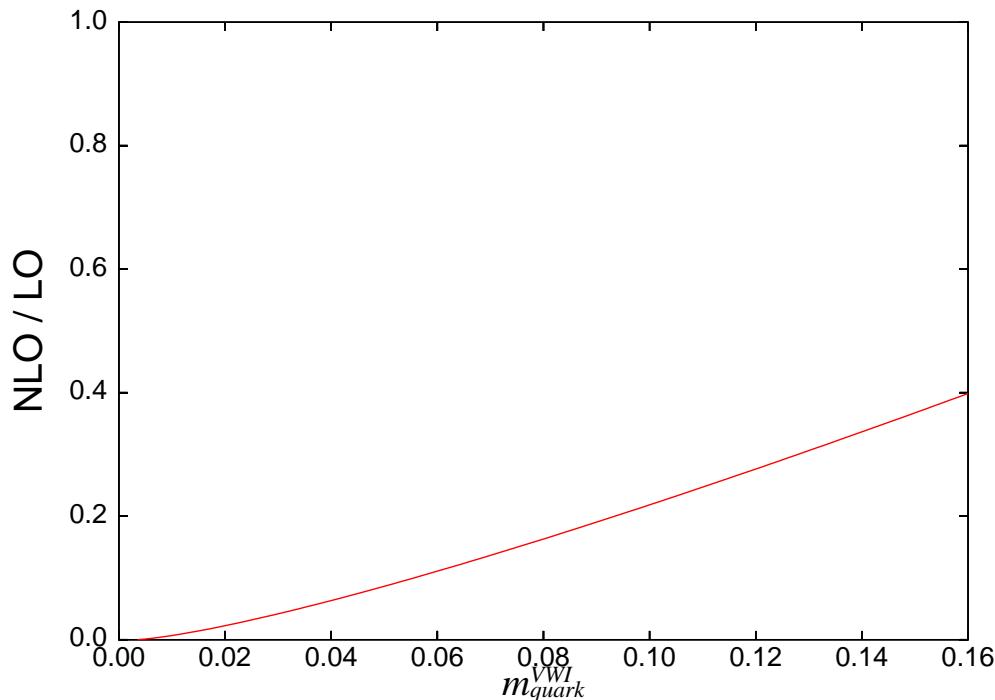
- The WChPT formulae describe our data well($\chi^2/dof = 1.6$), in contrast to the continuum ChPT case.
- $\Lambda_3 = 0.15(15)[\text{GeV}]$ is comparable with the phenomenological estimate($0.2 < \Lambda_3 < 2[\text{GeV}]$ [Gasser and Leutwyler,1984](#))
- $O(a)$ contribution ($\omega_1^{PS} - \omega_1^{AWI}$) suppresses the curvature in $m_{PS}^2/2m_{quark}^{AWI}$ to approximately 10% of the continuum ChPT

[Convergence check of the resummed WChPT]

Convergence is reasonable

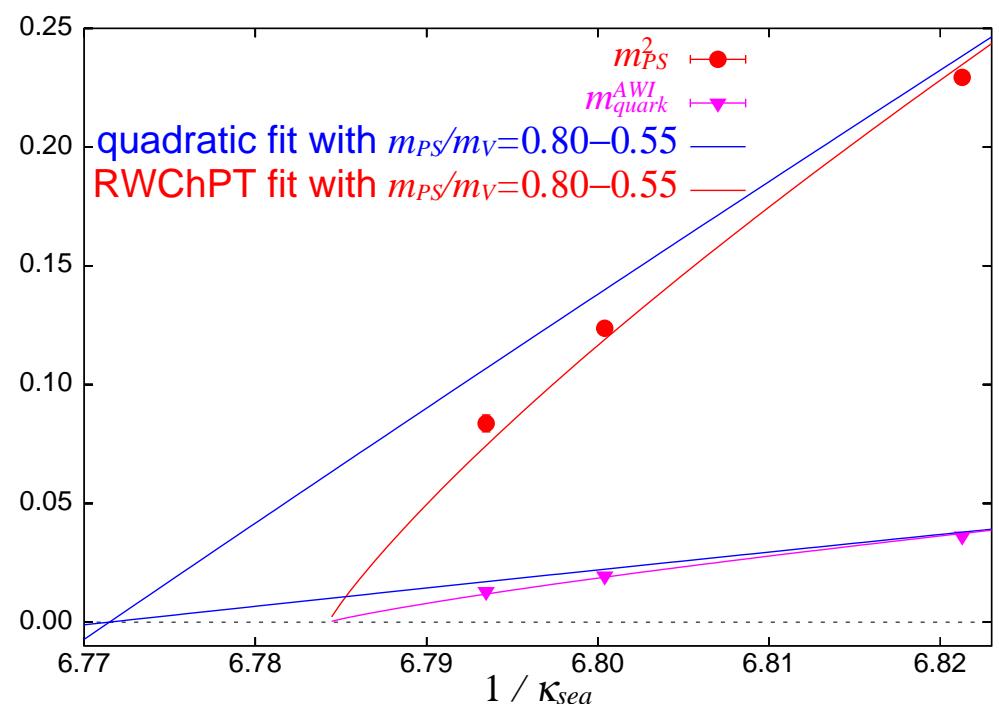
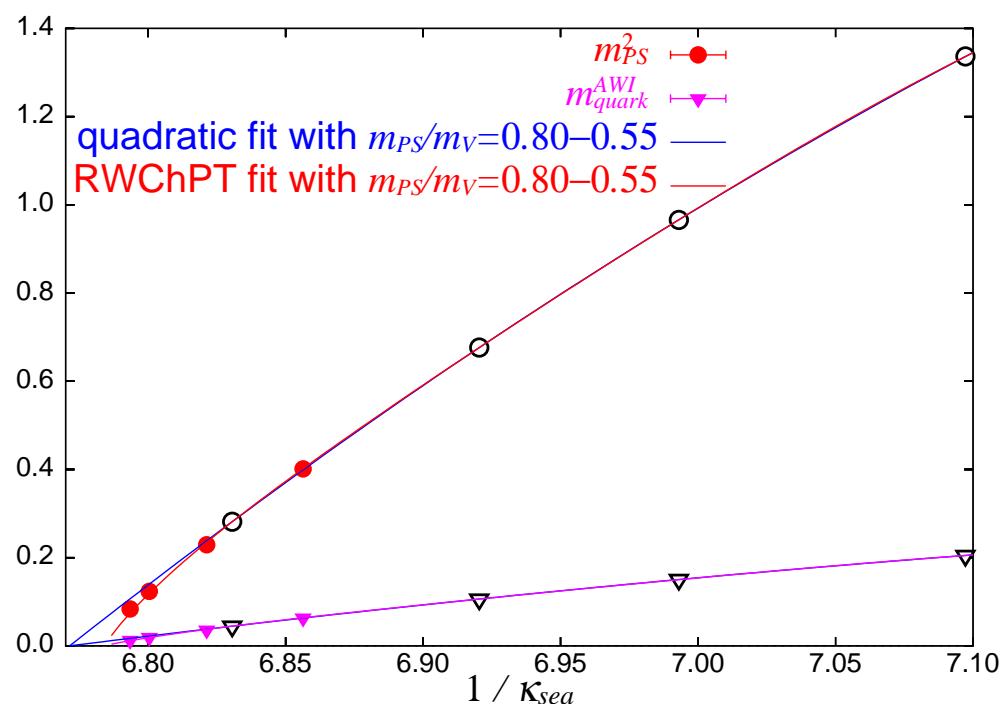
At the heaviest point,

- Tree : 71%
- one-loop : 29%



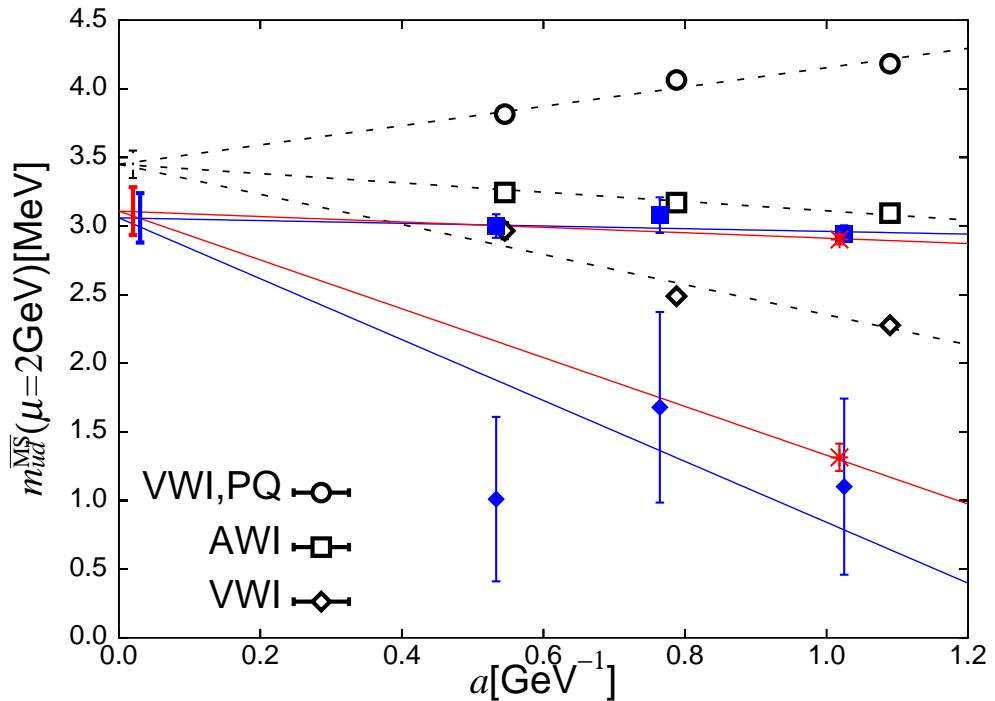
[Resummed WChPT extrapolation from intermediate quark masses]

- The WChPT formulae with the previous heavy data $m_{PS}/m_V = 0.80 - 0.55$ predicted the light data of $m_{PS}/m_V = 0.50 - 0.35$ better than polynomials
→ WChPT may work for intermediate quark masses



[Reanalysis of quark masses with the resummed WChPT]

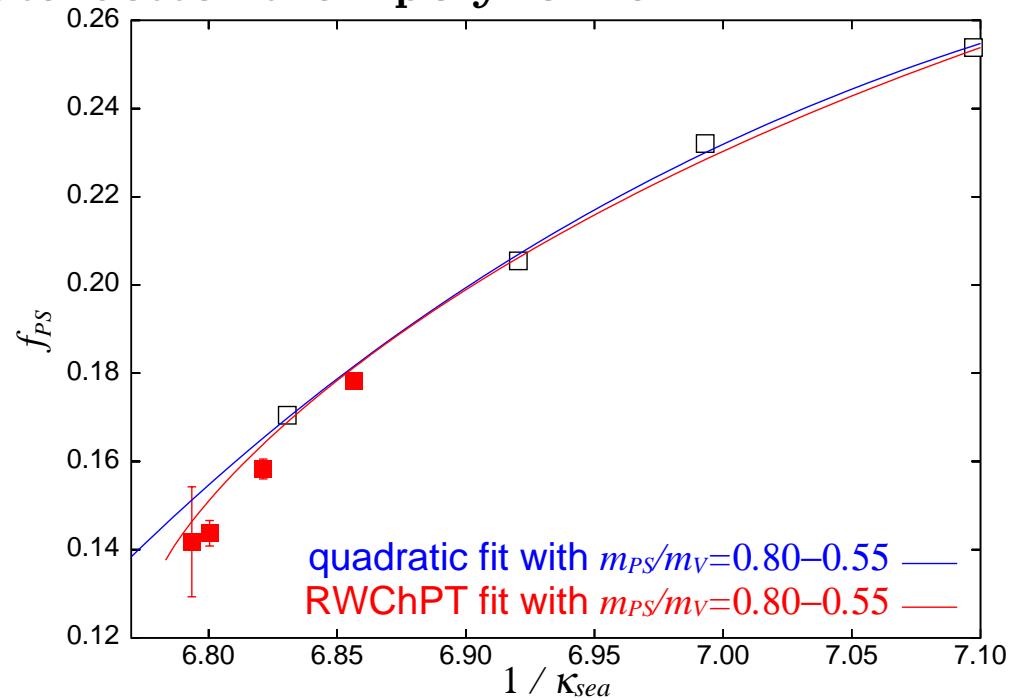
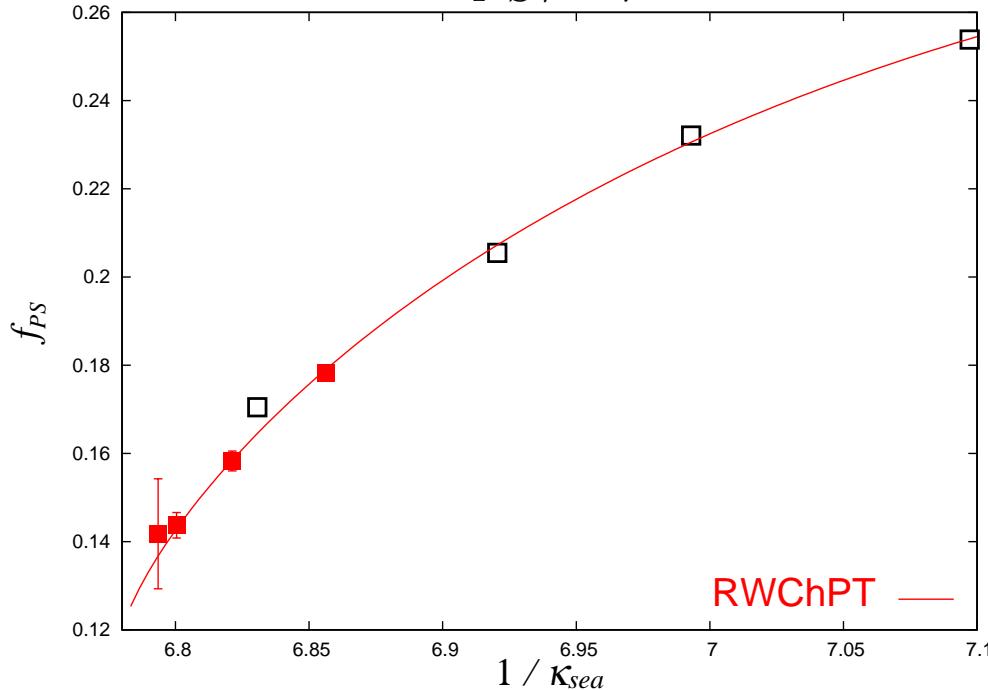
$$m_{ud}(\overline{\text{MS}}, \mu = 2\text{GeV}) = 3.11(17)\text{[MeV]} (3.45(10)\text{[MeV]}) \text{ cf. polynomial}$$



[WChPT formula for the pseudoscalar decay constant]

$$f_{PS} = f \left(1 - \omega_1^{f_{PS}} m_{quark}^{VWI} \log \left(\frac{A m_{quark}^{VWI}}{\Lambda_4^2} \right) \right)$$

- The WChPT describes our data of $m_{PS}/m_V = 0.80 - 0.35$ ($\chi^2/dof = 3.6$)
- $\Lambda_4 = 2.44(13)$ is comparable with(a little larger than) the phenomenological estimate ($\Lambda_4 = 1.26 \pm 0.14$ [GeV] [Colangelo et al.,2001](#))
→ The scaling check is needed
- WChPT formula with the previous data $m_{PS}/m_V = 0.80 - 0.55$
predicts $m_{PS}/m_V = 0.50 - 0.35$ data better than polynomial



5. Conclusion

We studied light hadron spectrum using several chiral extrapolations with our data of $m_{PS}/m_V = 0.80 - \textcolor{red}{0.35}$. Our results imply that **it is better to employ WChPT formulae instead of conventional polynomials for chiral extrapolations.**

[Polynomial extrapolation]

- Quadratic fits with heavy data at $m_{PS}/m_V = 0.80 - 0.55$ did not describe the behavior at small sea quark masses

[ChPT extrapolation]

- Our data do not show a quark mass dependence of ChPT
→ ChPT may be applicable only below $m_{PS}/m_V \sim 0.4$
cf. **UKQCD,qq+q 2004**

[WChPT extrapolation]

- WChPT ensures extrapolations toward small sea quark masses
- WChPT formulae describe our data of $m_{PS}/m_V = 0.80 - 0.35$ (in contrast to the continuum ChPT formulae)
- Convergence of one-loop WChPT formulae is reasonable
- WChPT fit to $m_{PS}/m_V = 0.80 - 0.55$ predicted $m_{PS}/m_V = 0.50 - 0.35$ data better than polynomials

[Prediction by WChPT formulae]

$$\begin{aligned} m_{ud}(\overline{\text{MS}}, \mu = 2\text{GeV}) &= 3.11(17) [\text{MeV}] \quad \text{by RWChPT} \\ \text{cf. } m_{ud}(\overline{\text{MS}}, \mu = 2\text{GeV}) &= 3.45(10) [\text{MeV}] \quad \text{by a polynomial} \end{aligned}$$

[Future works]

- Check the validity of WChPT formulae at smaller lattice spacings